

Nonconcave entropies from generalized canonical ensembles

Hugo Touchette

School of Mathematical Sciences

Queen Mary, University of London

London, UK

`ht@maths.qmul.ac.uk`

`http://www.maths.qmul.ac.uk/~ht`

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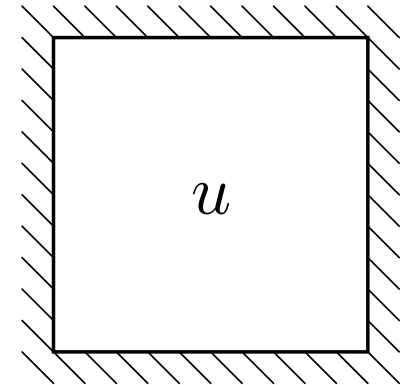


■ Context — plan

- Long-range systems have special properties
 - Their entropy can be nonconcave
- Nonconcavity breaks the Legendre transform of thermodynamics
 - Entropy \neq Legendre transform of canonical free energy
- How to calculate a nonconcave entropy?
 - Revision of concepts
 - Problem with nonconcave entropies
 - Generalized canonical ensemble
 - Example

■ Equilibrium statistical mechanics

- Microstate: $\omega \in \Omega$
- Hamiltonian: $U(\omega)$
- Mean energy: $u = U/n$



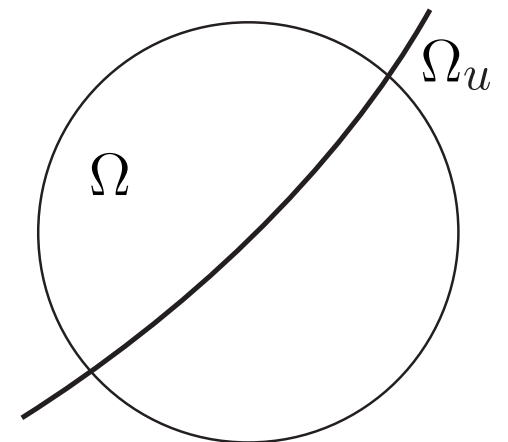
□ Microcanonical ensemble

- Density of states:

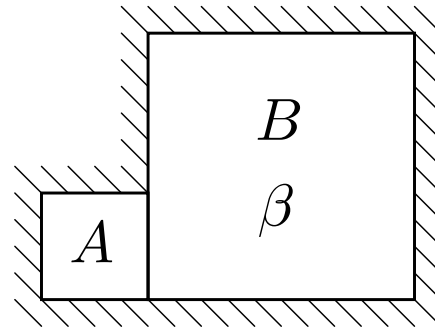
$$\begin{aligned}\rho(u) &= \int_{\Omega_u} d\omega \\ &= \int_{\Omega} \delta(u(\omega) - u) d\omega\end{aligned}$$

- Entropy function:

$$s(u) = \lim_{n \rightarrow \infty} \frac{1}{n} \log \rho(u)$$



□ Canonical ensemble



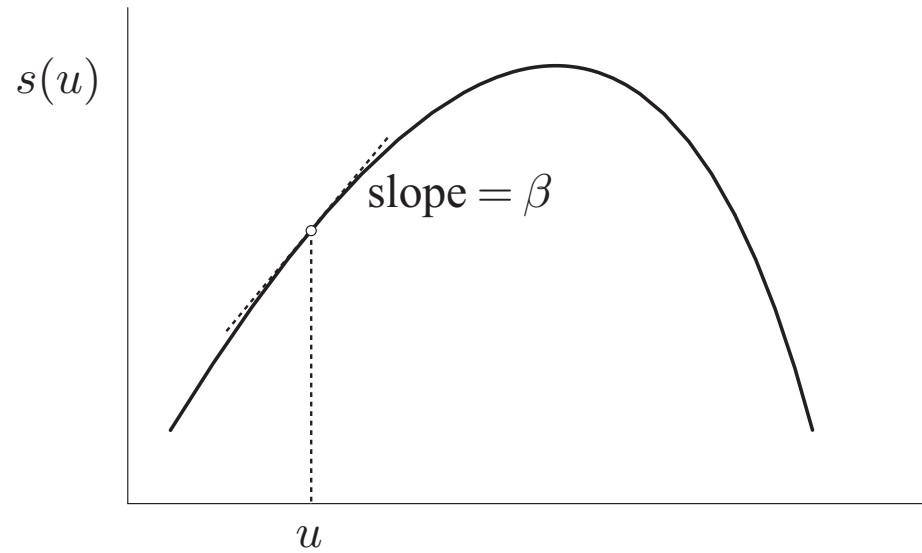
- Partition function:

$$Z(\beta) = \int_{\Omega} e^{-\beta U(\omega)} d\omega$$

- Free energy function:

$$\varphi(\beta) = - \lim_{n \rightarrow \infty} \frac{1}{n} \log Z(\beta)$$

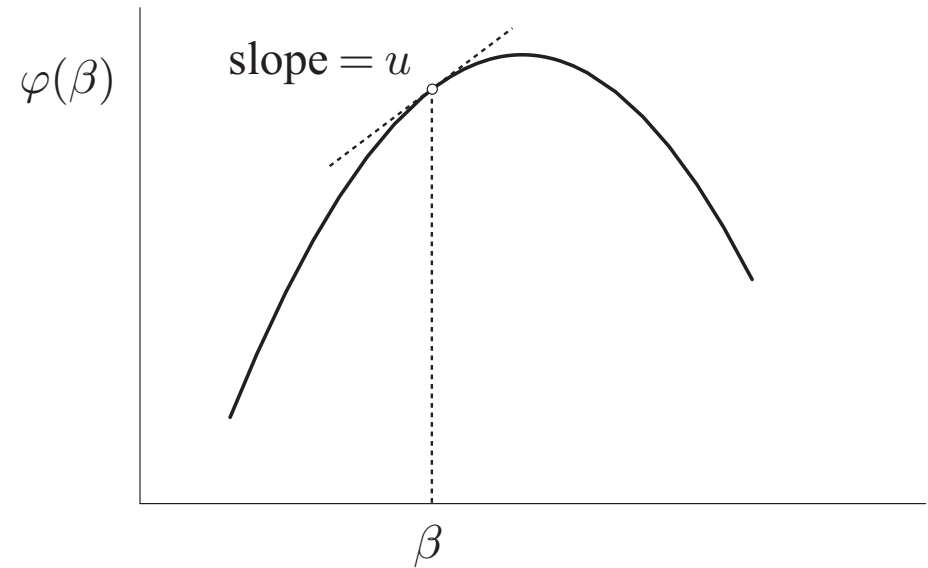
■ Calculation of the entropy function



$$s(u) = \beta u - \varphi(\beta)$$

$$\varphi'(\beta) = u$$

$$s = \varphi^*$$



$$\varphi(\beta) = \beta u - s(u)$$

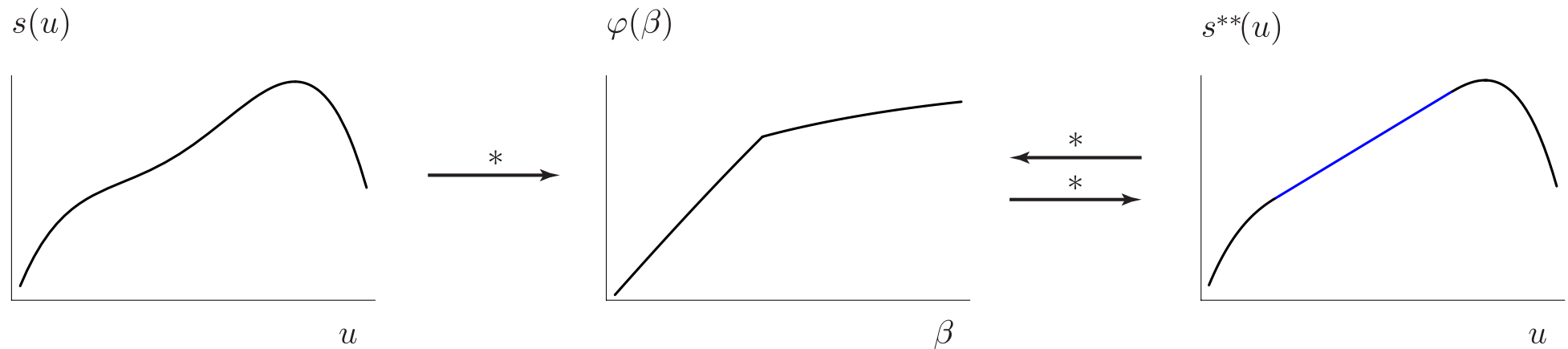
$$s'(u) = \beta$$

$$\varphi = s^*$$

$$s \longleftrightarrow \varphi$$

$$u \longleftrightarrow \beta$$

■ Problem with nonconcave entropies



Nonconcave

s

Always concave

$$\varphi = s^*$$

$$s \neq s^{**}$$

$$s^{**} = \varphi^*$$

- $s^{**}(u)$ is the concave envelope of $s(u)$
- $s^{**}(u) \geq s(u)$

■ Systems with nonconcave entropies

- Gravitating systems
- Spin models
 - Hamiltonian mean-field model
 - Blume-Emery-Griffiths model
 - Potts model
- 2D turbulence models
- String theory
- Multifractals
- Thermodynamic formalism of chaotic systems

■ Generalized canonical ensembles

Costeniuc, Ellis, Touchette & Turkington

JSP **119**, 1283 (2005); PRE **73**, 026105 (2006)

□ Definition

Standard canonical

$$P_\beta = \frac{e^{-n\beta u}}{Z(\beta)}$$

$$Z(\beta) = \sum_{\text{microstates}} e^{-n\beta u}$$

$$\varphi(\beta) = - \lim_{n \rightarrow \infty} \frac{1}{n} \log Z(\beta)$$

Generalized canonical

$$P_g = \frac{e^{-n\alpha u - ng(u)}}{Z_g(\alpha)}$$

$$Z_g(\alpha) = \sum_{\text{microstates}} e^{-n\alpha u - ng(u)}$$

$$\varphi_g(\alpha) = - \lim_{n \rightarrow \infty} \frac{1}{n} \log Z_g(\alpha)$$

□ Calculation of the entropy

- Concave entropy:

$$s = \varphi^*$$

- Nonconcave entropy:

$$s \neq \varphi^*$$

- Modified Legendre transform:

$$s = \varphi_g^* + g$$

Theorem

If, for a given g , $\varphi_g(\alpha)$ is differentiable at α , then

$$s(u) = \alpha u - \varphi_g(\alpha) + g(u), \quad \varphi_g'(\alpha) = u$$

■ Example: The 3-state mean-field Potts model

□ Definition of the model

- Microstate: $\omega = \omega_1, \omega_2, \dots, \omega_n$, $\omega_i \in \{1, 2, 3\}$

- Hamiltonian:

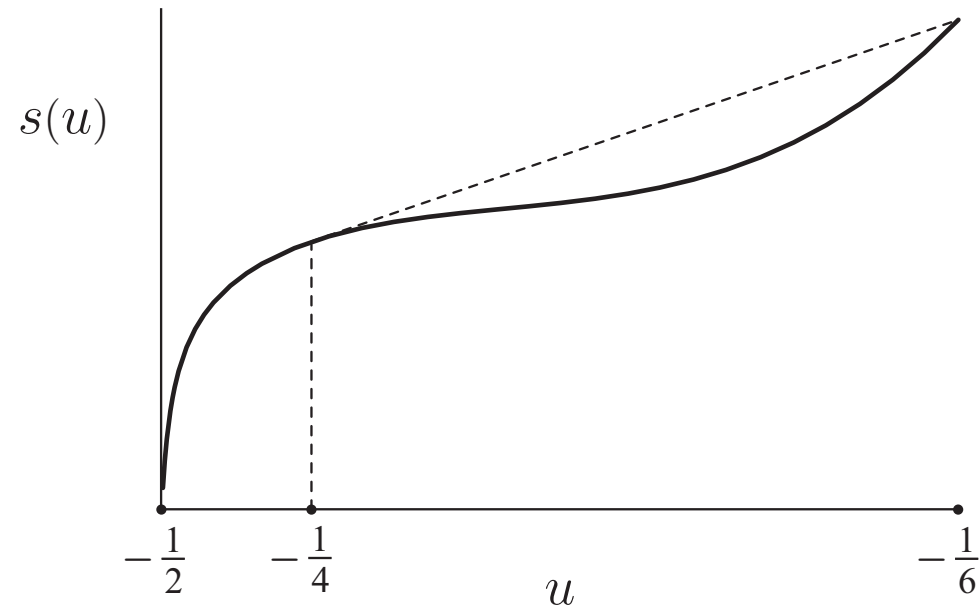
$$U(\omega) = -\frac{1}{2n} \sum_{i,k=1}^n \delta(\omega_i, \omega_k)$$

- Energy per spin:

$$u = \frac{U}{n} = -\frac{1}{2}(\nu_1^2 + \nu_2^2 + \nu_3^2), \quad \nu_i = \frac{\# \text{ spins } i}{n}$$

□ Entropy function

Costeniuc, Ellis & Touchette, J. Math. Phys. **46**, 063301, 2005



$$s(u) = -\frac{1 + \sqrt{2(-6u - 1)}}{3} \log\left(1 + \sqrt{2(-6u - 1)}\right) - \frac{2 - \sqrt{2(-6u - 1)}}{3} \log\left(\frac{2 - \sqrt{2(-6u - 1)}}{2}\right)$$

□ Gaussian ensemble

- g -function:

$$g(u) = \frac{\gamma}{2}u^2, \quad \gamma \in \mathbb{R}$$

- Partition function:

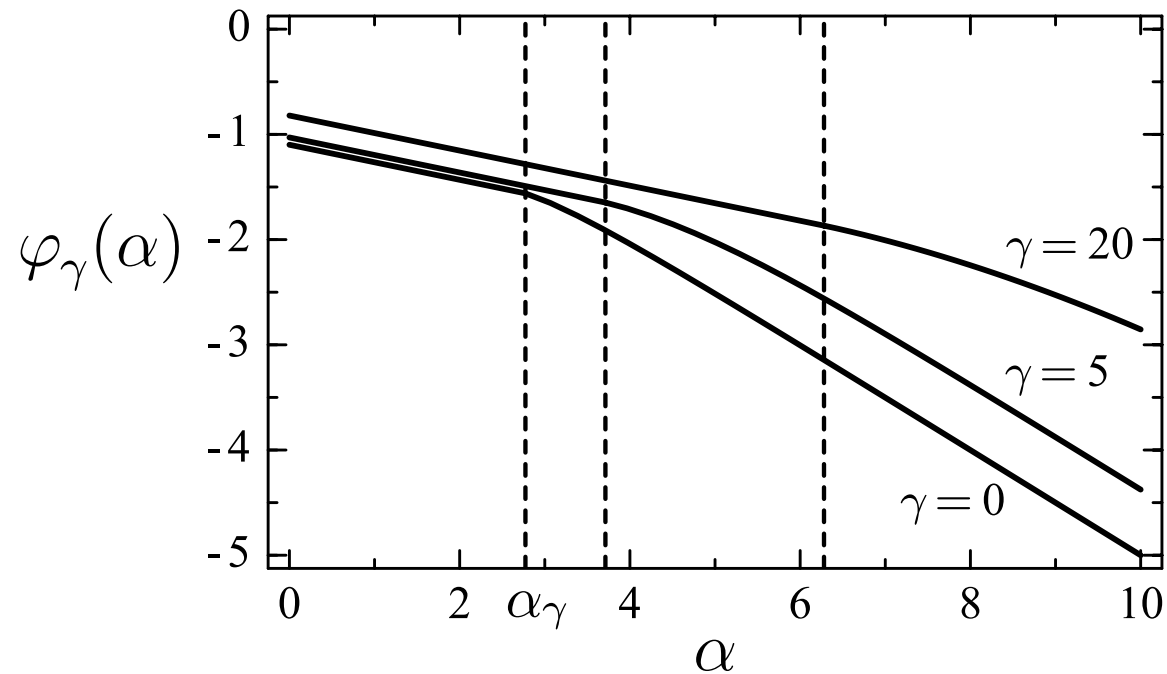
$$Z_\gamma(\alpha) = \sum_{\text{microstates}} e^{-n\alpha u - n\gamma u^2/2}, \quad \alpha, \gamma \in \mathbb{R}$$

- Free energy:

$$\varphi_\gamma(\alpha) = - \lim_{n \rightarrow \infty} \frac{1}{n} \log Z_\gamma(\alpha)$$

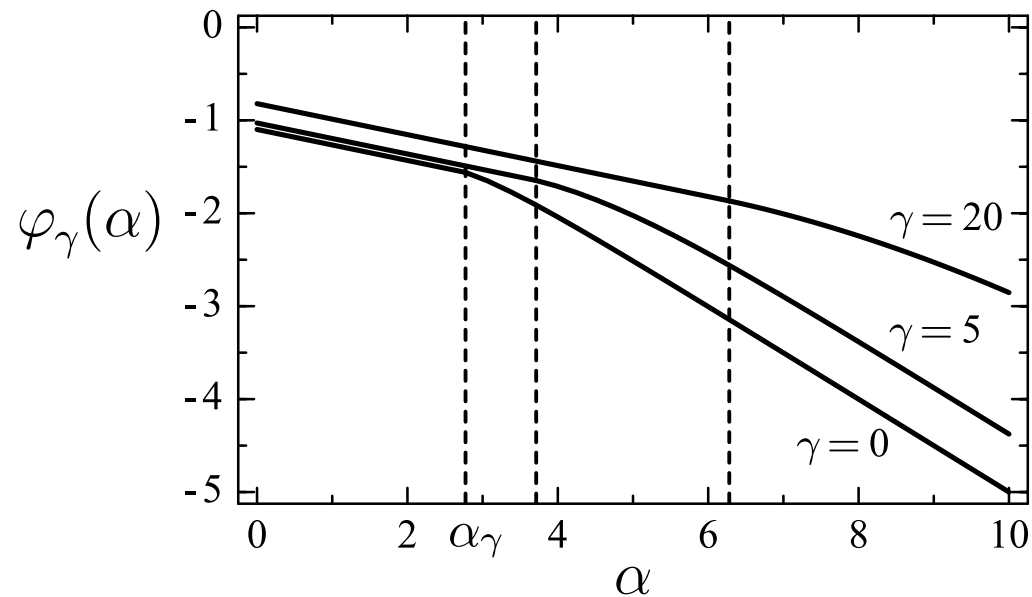
□ Calculation of the Gaussian free energy

Costeniuc, Ellis & Touchette, cond-mat/0605213, to appear in PRE

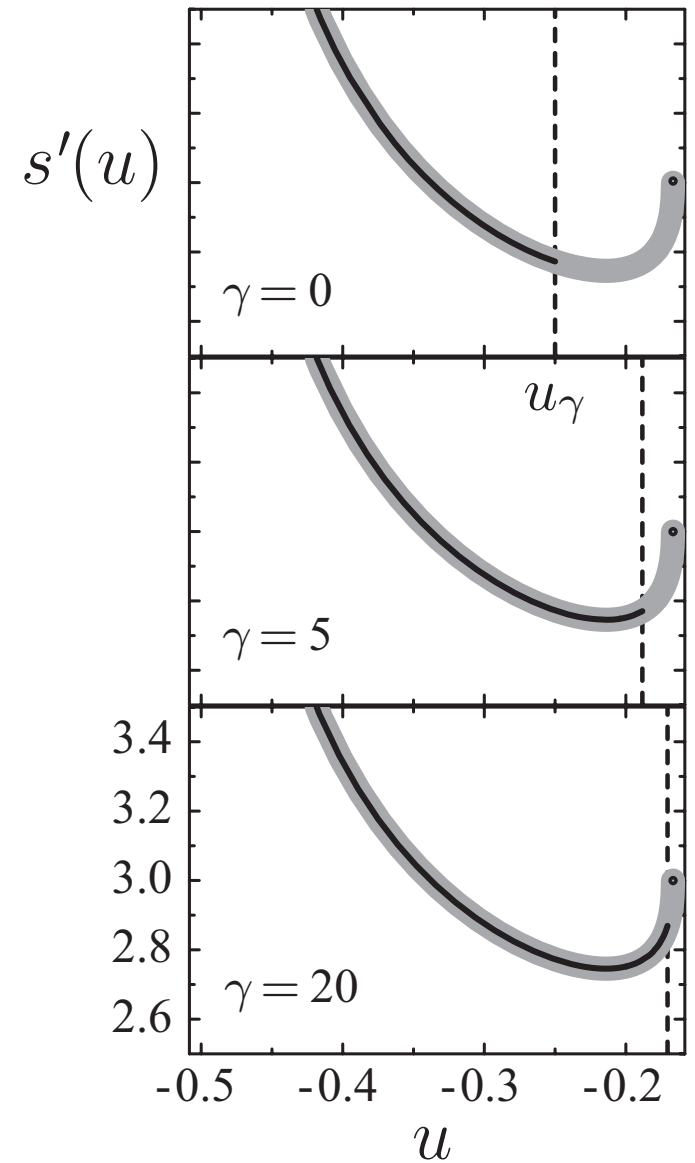


- Non-differentiable point: α_γ
- Left branch: $\varphi'_\gamma(\alpha) = -\frac{1}{6}$
- Right branch: $\varphi'_\gamma(\alpha) \in (-\frac{1}{2}, u_\gamma)$

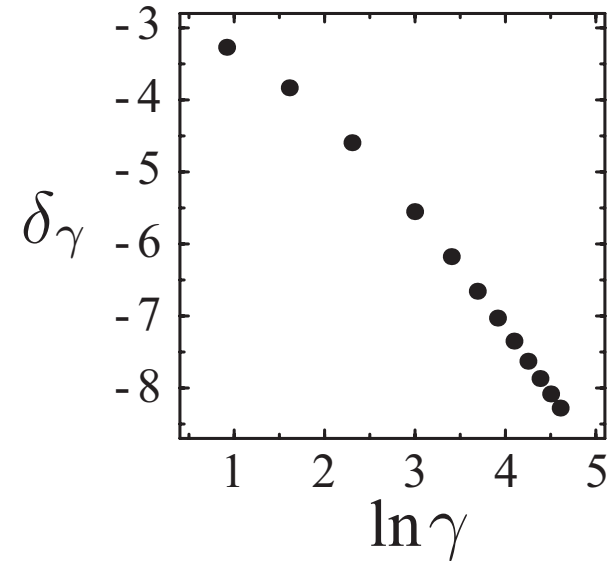
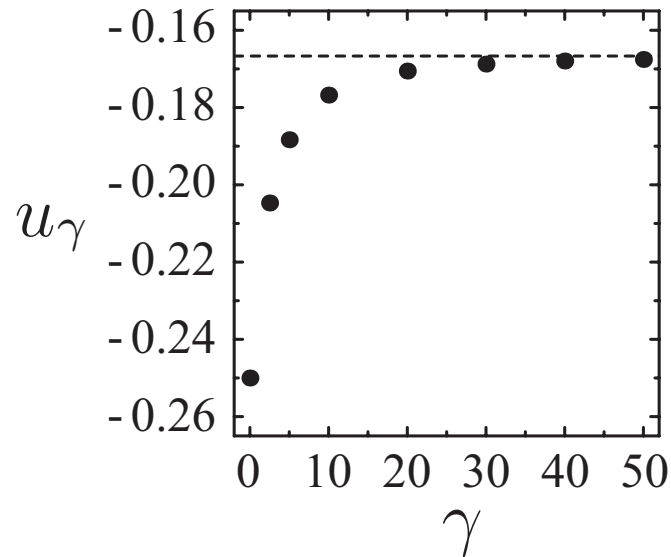
□ Calculation of the entropy



- $s = \varphi_g^* + g$
- $s(u) = \alpha u + \frac{\gamma}{2}u^2 - \varphi_\gamma(\alpha)$
- $u = \varphi'_\gamma(\alpha) \in (-1/2, u_\gamma)$



□ Asymptotic equivalence



- Maximum mean energy: $u_\gamma = \varphi'_\gamma(\alpha_\gamma + 0)$
- Asymptotic limit:

$$u_\gamma \rightarrow u_{\max} = -\frac{1}{6} \quad \text{as} \quad \gamma \rightarrow \infty$$

- Scaling:

$$\delta_\gamma = \ln |u_\gamma - u_{\max}| \sim -2 \ln \gamma$$

□ Equivalence at the macrostate level

- Macrostate:

$$a = \frac{\# \text{ spin } 1}{n}$$

- Gaussian equilibrium:

$$a_{\gamma, \alpha}$$

- Microcanonical equilibrium:

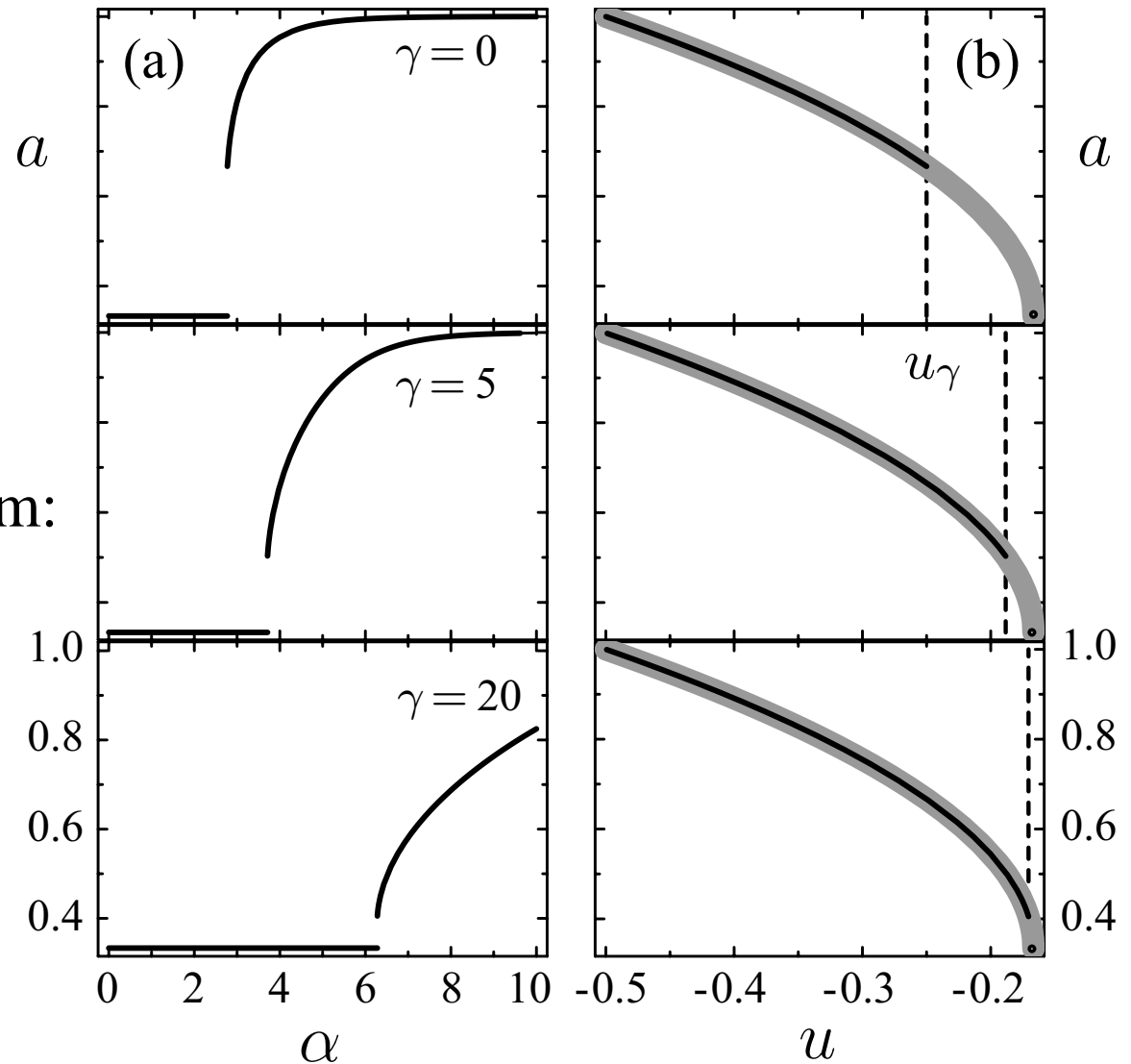
$$a^u$$

- $a^{-\frac{1}{2}} = 1$

- $a^{-\frac{1}{6}} = \frac{1}{3}$

- Mapping:

$$a \rightarrow (u(a), a)$$



■ Summary

- Entropy can be **nonconcave** (for long-range systems)
- Related to first-order phase transitions
- Nonconcave entropies cannot be calculated in the canonical ensemble

$$s \neq \varphi^*$$

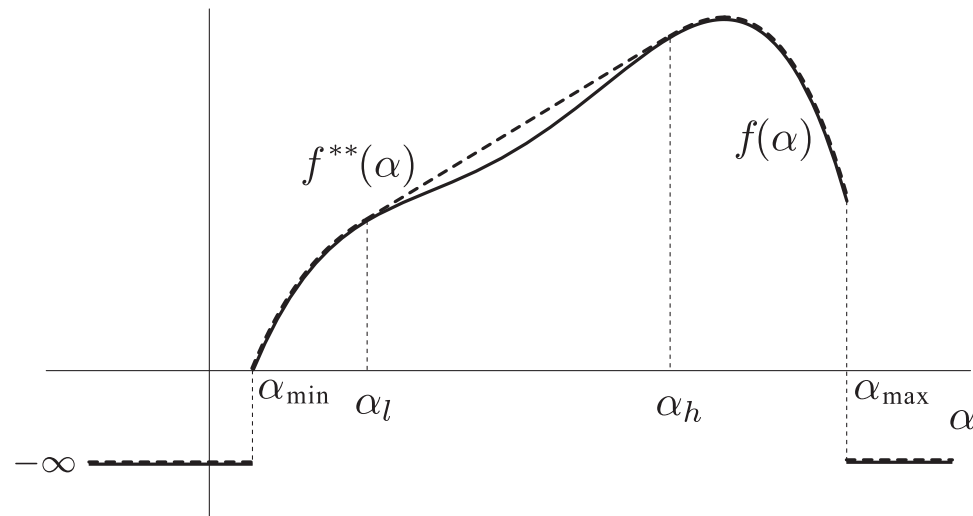
- Fix the problem: generalized canonical ensembles

$$s = \varphi_g^* + g$$

- Gaussian ensemble: $g(u) = \gamma u^2$
 - **Any** shape of entropy can be obtained via the Gaussian ensemble
 - Gaussian ensemble is **universal**
- Other ensembles possible
 - $g(u) = \gamma|u|$, $g(u) = \gamma u^4$, ...

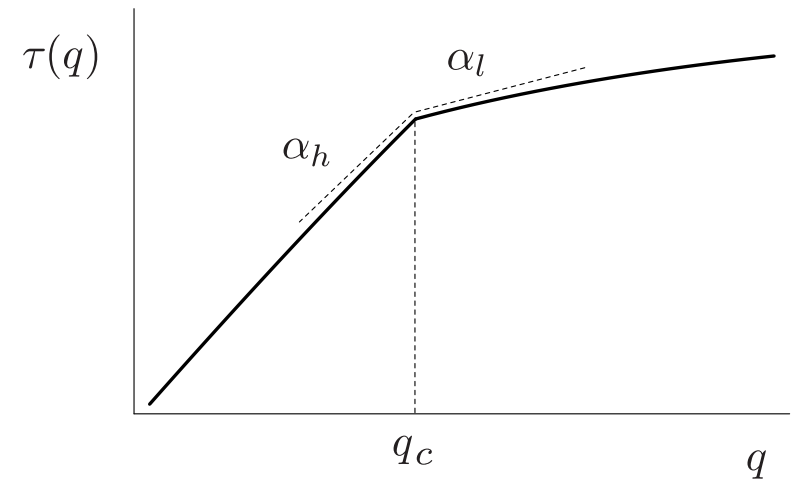
■ Nonconcave multifractals

Touchette & Beck, cond-mat/0507379, to appear in JSP



$$f \neq \tau^*$$

$$f = \tau_g^* + g$$



$$\tau = f^*$$

■ Other examples

- Thermodynamic formalism

- Nonconvex spectra of dynamical indices:

$$P(\lambda) \sim e^{-t\psi(\lambda)}$$

- Fluctuation theorems

- Nonconvex fluctuation functions:

$$P(W_t/t = w) \sim e^{-tf(w)}$$

- Large deviation theory

- Nonconvex rate functions:

$$P(A_n = a) \sim e^{-nf(a)}$$

■ References

- M. Costeniuc, R.S. Ellis, H. Touchette, [cond-mat/0605213](#), to appear in PRE
- H. Touchette, C. Beck, [cond-mat/0507379](#), to appear in JSP
- M. Costeniuc, R.S. Ellis, H. Touchette, B. Turkington, *J. Stat. Phys.* **119**, 1283-1329, 2005.
- M. Costeniuc, R.S. Ellis, H. Touchette, B. Turkington, *Phys. Rev. E* **73**, 026105, 2006.

Work in collaboration with

Richard S. Ellis (UMass, USA)

Marius Costeniuc (Max Planck Inst., Germany)

■ Everything is extensive!

- Energy:

$$U \sim n, \quad u = \frac{U}{n}$$

- Entropy:

$$\ln |\Omega_u| \sim n$$

$$s(u) = \lim_{n \rightarrow \infty} \frac{1}{n} \ln |\Omega_u|$$

- Free energy:

$$\ln Z(\beta) \sim n$$

$$\varphi(\beta) = \lim_{n \rightarrow \infty} -\frac{1}{n} \ln Z(\beta)$$